					Table 2	(cont.)					
hkl	$oldsymbol{F}_{ ext{obs.}}$	$oldsymbol{F}_{ ext{calc.}}$	α (°)	hkl	$F_{ m obs.}$	${F}_{ m calc.}$	α (°)	hkl	$oldsymbol{F}_{ ext{obs.}}$	${F}_{ m calc.}$	α (°)
512	8.0	8.3	344	6.11.0	5.9	5.9	270	720	0.9	0.2	0
522	5.9	5.6	26					730	$2 \cdot 1$	$2 \cdot \overline{2}$	270
532	5.3	3.4	123	601	0.9	0.2	270	740	3.8	3.0	180
$\bf 542$	5.7	5.0	338	611	$2 \cdot 3$	1.2	197	750		1.3	90
552	6.5	5.7	338	621	$5 \cdot 3$	3.1	80	760	1.4	0.2	0
562	3.8	4.3	107	631	3.0	$2 \cdot 3$	112	770		0.7	270
572	4.4	$4 \cdot 3$	285	641	4.9	5.7	355	780	1.8	0.7	0
582	$2 \cdot 7$	$2 \cdot 9$	299	651	$2 \cdot 0$	2.5	15				
$\bf 592$	$2 \cdot 2$	$2 \cdot 8$	88	661	$5 \cdot 1$	$3 \cdot 9$	86	701	$5 \cdot 1$	6.7	90
5.10.2	1.5	1.9	101	671	$2 \cdot 6$	$2 \cdot 4$	138	711	3.8	3.6	316
				681	$2 \cdot 0$	$2 \cdot 6$	355	721	6.7	6.4	259
503		1.1	270	691	4.3	3.6	348	731	3.9	3.9	6
513	$6 \cdot 4$	4.6	205	6.10.1	1.7	$2 \cdot 3$	113	741	4.0	3.5	213
523	$2 \cdot 6$	3.6	264					751	$2 \cdot 7$	$2 \cdot 1$	316
533	11.7	12.3	171	602	$9 \cdot 2$	10.6	180	761	1.1	$2 \cdot 3$	114
543	$2 \cdot 9$	$3 \cdot 4$	42	612	$9 \cdot 0$	8.5	237	771	1.7	$1 \cdot 2$	269
553		$2 \cdot 3$	348	622	1.9	$2 \cdot 3$	17	781	3.5	4.5	249
563	$3 \cdot 2$	$3 \cdot 4$	264	632	$2 \cdot 9$	1.7	327				
				642	$3 \cdot 2$	3.0	257	702	0.7	0.8	180
504		0.4	180	652	6.7	$7 \cdot 0$	233	712	$2 \cdot 6$	1.1	168
514	$2 \cdot 7$	$2 \cdot 1$	238	662	3.9	3.8	161	722	1.3	1.3	6
$\bf 524$	$3 \cdot 4$	$3 \cdot 2$	338	672	3.7	3.9	264	732	3.4	3.1	26
534	$2 \cdot 9$	$2 \cdot 4$	120	682	$2 \cdot 7$	4.1	332	$\bf 742$	1.7	$2 \cdot 0$	90
544		0.7	23				ŀ				
				603	$5 \cdot 0$	$5 \cdot 1$	90	703	1.0	0.7	270
600	$2 \cdot 7$	1.4	180	613	$2 \cdot 7$	$2 \cdot 7$	302	713	$1 \cdot 3$	1.0	162
610	7.7	$6 \cdot 6$	270	623	0.8	0.9	201				
620	$5 \cdot 1$	3.9	0	633	4.1	$4 \cdot 2$	17	800	6.5	4.3	0
630	$1 \cdot 2$	$1 \cdot 2$	90	643	$2 \cdot 9$	2.8	103	810	1.1	0.8	90
640	1.0	$3 \cdot 0$	0	653	$2 \cdot 2$	$2 \cdot 6$	205	820	$1 \cdot 4$	1.9	0
650	7.5	$8\cdot 2$	270					830	$2 \cdot 1$	$2 \cdot 4$	90
660	$5 \cdot 2$	6.8	180	604	$2 \cdot 9$	$3 \cdot 4$	180				
670	6.7	5.7	270	614	$4 \cdot 2$	4.0	303	801	$3 \cdot 0$	$2 \cdot 3$	90
680	$2 \cdot 9$	1.6	180					811	$3 \cdot 2$	$2 \cdot 8$	88
690	1.1	0.7	270	700	8.1	$8 \cdot 3$	180	821	3.9	$1 \cdot 3$	299
6.10.0		0.1	180	710	$2 \cdot 5$	$1 \cdot 2$	270				

Short Communications

Contributions intended for publication under this heading should be expressly so marked; they should not exceed about 500 words; they should be forwarded in the usual way to the appropriate Co-editor; they will be published as speedily as possible; and proofs will not generally be submitted to authors. Publication will be quicker if the contributions are without illustrations.

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A new punched-card method of Fourier synthesis. By D. M. S. Greenhalgh and G. A. Jeffrey, Department of Inorganic and Physical Chemistry, University of Leeds, England

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In the Hollerith methods for summing Fourier series which have been described hitherto (e.g. Cox, Gross & Jeffrey, 1947) the punched cards have been essentially more accurate Beevers-Lipson strips. It was possible to reduce the number of cards by using the technique of 'progressive digiting', and subsequently we found that further advantages could be obtained by expressing the amplitudes in the binary scale and using the feature of the B.T.M. Rolling Total Tabulator whereby a counter can be added into itself to double its contents; nevertheless, all methods retained an undesirable feature of the original Beevers-Lipson procedure, namely, the preliminary selection of cards from a filing system. Hand selection requires careful checking, and ultimately the cards must be refiled and again checked; these are tedious and time-consuming operations, however well they are organized.

This note outlines the principles of a new method which has been developed by one of us (D. M. S. Greenhalgh) and which has the great advantage that hand selection and filing of cards have been eliminated. Full technical details are being published elsewhere (Greenhalgh, 1950). For the first stage of a synthesis, the appropriate amplitudes A_n are hand-punched on to the relatively small number of cards required, and for subsequent stages the amplitudes can be summary-punched from the results of the preceding summations. The trigonometric functions are represented in the binary scale in the way described below, and are transferred to these cards from a master pack. The whole set of cards can be stored afterwards as a compact and permanent record of the job.

All the arithmetical operations in a synthesis by the new method are performed in the ordinary decimal scale on the

amplitudes, the binary digits (0 or 1) of the trigonometrical function (referred to below as the multiplier) serving only to give instructions to the tabulator. Fourteen binary digits are required to express one value of the function $10,000\cos 2\pi nx$ (or $10,000\sin 2\pi nx$), and each digit is represented on a card by a hole for 1 or a blank for 0. Only one punching position is required to represent one binary digit (so that since one column of 12 punching positions is required to represent one decimal digit for conventional arithmetical operations, about four (~12/log₂10) times as much information can be carried on a card by the new method). Since the 14 digits of one particular multiplier are required to give 14 instructions in succession to the tabulator they are represented by the same punching position on 14 separate cards. The values of cos 2nx for a fixed value of n and values of x from 0 to $\frac{1}{2}$ at intervals of $\frac{1}{120}$, for example, are represented by 31 punching positions (less than 3 columns) on 14 cards. There is one such pack of 14 cards for each order n, and the appropriate value of the (three-figure) A_n is punched in the normal (decimal) way in the same columns on all 14 cards (negative amplitudes being punched as complements in four columns).

To carry out the summation $\sum\limits_{0}^{N}A_{n}\cos2\pi nx$ all the

14(N+1) cards are sorted into groups in the order of the binary digits. They are then tabulated in such a way that the amplitude is fed into the counter only if the corresponding binary digit on the card is 1. At the change of the binary digit, i.e. from the first to the second and so on, the contents of the counter are added to itself. Suppose in a very simple synthesis the data are as follows (the missing A_n 's being all zero):

		Multiplier (e.g. $10 \sin 2nx_r$)						
\boldsymbol{n}	A_n	Decimal		Binary				
1	120	2	0	0	1	0		
4	65	7	0	1	1	1		
6	20	. 9	1	0	0	1		
8	10	10	1	0	1	0		
			\boldsymbol{a}	b	c	d		

The following tabulator operations are required to obtain $\Sigma A_n \sin 2\pi n x_r = 97.5$:

(a) Add $20 + 10$ into the counter and double	60
(b) Add 65 and double again	250
(c) Add $10+65+120$ and double again	890
(d) Add $20 + 65$	975

In practice negative multipliers occur, and these may be dealt with either

 (i) by adding 1 to all the sine and cosine values before converting to the binary scale, and subsequently making the necessary subtraction;

or (better)

(ii) by operating in the 'scale of minus two'. Any number, positive or negative, may be split into powers of -2, and the only change needed to multiply in this scale is to punch the amplitudes for alternate digits with signs reversed.

Since the data for each one-dimensional synthesis occupy only a few columns of the cards, the information for as many syntheses as convenient are punched into one pack of cards: for example, with 3-figure amplitudes and with syntheses extending over one-quarter of a cell side at $\frac{1}{120}$ th intervals there is space for 18 one-dimensional syntheses. It is then possible, for example, at one passage of the same cards, to calculate the electron density at a number of points simultaneously either along the line (x, y_1, z_1) or along the line (x_r, y, z_1) . This flexibility is of considerable value in exploring limited portions of an electron-density distribution.

We are indebted to Prof. E. G. Cox for providing facilities for this work and for his interest in it.

References

Cox, E. G., Gross, L. & Jeffrey, G. A. (1947). Proc. Leeds Phil. Soc. 5, 13.Greenhalgh, D. M. S. (1950). Proc. Leeds Phil. Soc.

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Anomalous X-ray reflexions from copper crystals. By W. A. Rachinger, Baillieu Laboratory, University of Melbourne, Australia (Received 28 December 1949)

In a recent investigation of the plastic properties of copper single crystals, it was found that when a deformed crystal was set in position for the high-angle (004) Co $K\alpha$ reflexion there sometimes appeared another line close to the main (004) reflexion. The relative position of this line and the (004) reflexion is shown in Fig. 1, which is a flat-film back-reflexion diagram. The extraneous line has a doublet structure and is not concentric with the (004) Debye circle. This latter fact rules out the possibility of the line being a primary Bragg reflexion due to radiation of another wave-length. It was also found that the reflexion was not a 'slit-effect' or due in any way to the experimental arrangement.

The line can, in fact, be shown to be due to a double reflexion of the incident X-ray beam, first by a $(1\overline{1}1)$ and then by a $(1\overline{1}3)$ plane. This is illustrated by Fig. 2, which represents the trace of the (110) plane across the standard stereographic projection of a cubic crystal. The $(1\overline{1}1)$,

(001) and (113) poles are represented, not as points, but as circular areas since, as will be seen later, the double reflexion will accompany the primary (004) reflexion only in crystals which are imperfect in that their reflecting planes have a considerable orientation range. If the crystal is set so that the X-ray beam lies parallel to the (110) plane and is inclined at 8.3° to the mean position of the (001) normals, then the direction of the beam will be represented by the point A in the standard projection. Now with the crystal in this position the (004) reflexion, for which $\theta = 81.7^{\circ}$, will take place and the direction of the reflected beam will be represented by the point B. The (111) reflexion ($\theta = 25.4^{\circ}$) will also take place provided that the range of the (111) normals is greater than $\pm 1.6^{\circ}$, since a point A' can then be found in the $(1\overline{1}1)$ region which is $90 - \theta (= 64.6^{\circ})$ from the incident beam A. The ray reflected from this plane A' will be represented by the point C. This (111) reflexion can be further reflected by